# PHYS 234: Quantum Physics 1 (Fall 2008) <br> Assignment 9 - Solutions 

1. Two possible eigenfunctions for a particle moving freely in a region of length $a$, but strictly confined to that region, are shown in the figure below. When the particle is in the state corresponding to the eigenfunction $\psi_{\mathrm{I}}$, its total energy is 4 eV .

a) What is its total energy in the state corresponding to $\psi_{\text {II }}$ ?
b) What is the lowest possible energy for the particle in this system?
solution a) In the lowest energy state $n=1, \psi$ has no nodes. Hence $\psi_{I}$ must correspond to $n=2$, and $\psi_{I I}$ to $n=3$. Since the energy of the $n^{t h}$ state $E_{n} \propto n^{2}$ and $E_{I}=4 \mathrm{eV}$, then

$$
\frac{E_{I I}}{E_{I}}=\frac{3^{2}}{2^{2}} ; \quad E_{I I}=9 \mathrm{eV}
$$

b) By the same analysis,

$$
\frac{E_{0}}{E_{I}}=\frac{1^{2}}{2^{2}} ; \quad \quad E_{I I}=1 \mathrm{eV}
$$

2. It can be proved that, in general, the inner product of any two different eigenfunctions (of the same given potential) is always zero. This property is called orthogonality. The inner product is defined as,

$$
\int_{-\infty}^{\infty} \psi_{m}(x)^{*} \psi_{n}(x) d x=0 \quad m \neq n
$$

Form the inner product of the eigenfunctions for the $n=1$ and $n=3$ states of the infinite square well potential. Show that the inner product is equal to zero. In other words, show that

$$
\int_{-\infty}^{\infty} \psi_{1}^{*}(x) \psi_{3}(x) d x=0
$$

Hint: The following identities may be useful:

$$
\begin{aligned}
\cos (v) \cos (v) & =[\cos (u+v)+\cos (u-v)] / 2) \\
\cos (u)-\cos (v) & =-2 \sin \left(\frac{u+v}{2}\right) \sin \left(\frac{u-v}{2}\right)
\end{aligned}
$$

solution The inner product that is formed is analysed as follows

$$
\int_{-\infty}^{+\infty} \psi_{1}^{*} \psi_{3} d x=\frac{2}{a} \int_{-a / 2}^{+a / 2} \cos \left(\frac{\pi x}{a}\right) \cos \left(\frac{3 \pi x}{a}\right) d x
$$

The integral is simplified using the substitution $u=\frac{\pi x}{a}$ and $v=\frac{3 \pi x}{a}$ and the relationship from the problem sheet $\cos (v) \cos (v)=[\cos (u+v)+\cos (u-v)] / 2$, hence ${ }^{a}$

$$
\int_{-\infty}^{+\infty} \psi_{1}^{*} \psi_{3} d x=\frac{1}{a} \int_{-a / 2}^{+a / 2} \cos \left(\frac{4 \pi x}{a}\right)+\cos \left(\frac{2 \pi x}{a}\right) d x
$$

Both terms in the integral integrate to zero over the range of integration, hence

$$
\int_{-\infty}^{+\infty} \psi_{1}^{*} \psi_{3} d x=0
$$

which is the orthogonality property which we needed to prove.
3. The wave function for a particle is

$$
\Psi(x, t)=\sin (k x)[i \cos (\omega t / 2)+\sin (\omega t / 2)]
$$

where $k$ and $\omega$ are constants.
a) Is this particle in a state of definite momentum? If so, determine the momentum.
b) Is this particle in a state of definite energy ? If so, determine the energy.
solution A wave function representing a particle with a definite value for an observable is an eigenfunction of the operator for that observable, with the eigenvalue being the value of that observable and hence the result of a measurement.
a) Apply the momentum operator:

$$
-i \hbar \frac{\partial}{\partial x} \sin (k x)[i \cos (\omega t / 2)+\sin (\omega t / 2)]=i \hbar k \cos (k x)[i \cos (\omega t / 2)+\sin (\omega t / 2)]
$$

This is not a constant times the original wave function, and so not an eigenfunction of the momentum operator and therefore the particle is not in a state with well-defined momentum.
b) Apply the energy operator:

$$
\begin{aligned}
i \hbar \frac{\partial}{\partial t} \sin (k x)[i \cos (\omega t / 2)+\sin (\omega t / 2)] & =i \hbar \sin (k x)[-i(\omega / 2) \sin (\omega t / 2)+(\omega / 2) \cos (\omega t / 2)] \\
& =(\hbar \omega / 2) \sin (k x)[i \cos (\omega t / 2)+\sin (\omega t / 2)]
\end{aligned}
$$

This is a constant times the original wave function, and so the wave function is an eigenfunction of the energy operator, and therefore the particle is in a state of definite energy. The value of the energy is the eigenvalue, $E=\hbar \omega / 2$
4. Using the first two normalized wave functions $\Psi_{1}(x, t)$ and $\Psi_{2}(x, t)$ for a particle moving freely in a region of length $a$, but strictly confined to that region, construct the linear combination

$$
\Psi(x, t)=c_{1} \Psi_{1}(x, t)+c_{2} \Psi_{2}(x, t)
$$

which is a superposition of the first two energy eigenstates. Then derive a relation involving the adjustable constants $c_{1}$ and $c_{2}$ which, when satisfied, will ensure that $\Psi(x, t)$ is also normalised.
solution The wavefunctions in question are

$$
\Psi_{1}=\sqrt{\frac{2}{a}} \cos (\pi x / a) \exp \left(-i E_{1} t / \hbar\right) ; \Psi_{2}=\sqrt{\frac{2}{a}} \sin (2 \pi x / a) \exp \left(-i E_{2} t / \hbar\right)
$$

with $E_{2}=4 E_{2}$. The linear combination is

$$
\Psi=c_{1} \Psi_{1}+c_{2} \Psi_{2}
$$

Normalising this gives

$$
1=\int_{-\infty}^{\infty} \Psi^{*} \Psi d x
$$

Substituting and expanding, this becomes

$$
c_{1} c_{1}^{*} \int_{-\infty}^{\infty} \Psi_{1}^{*} \Psi_{1} d x+c_{2} c_{2}^{*} \int_{-\infty}^{\infty} \Psi_{2}^{*} \Psi_{2} d x+c_{1}^{*} c_{2} \int_{-\infty}^{\infty} \Psi_{1}^{*} \Psi_{2} d x+c_{2}^{*} c_{1} \int_{-\infty}^{\infty} \Psi_{2}^{*} \Psi_{1} d x=1
$$

Since $\Psi_{1}$ and $\Psi_{2}$ are already normalised,

$$
\int_{-\infty}^{\infty} \Psi_{1}^{*} \Psi_{1} d x=\int_{-\infty}^{\infty} \Psi_{2}^{*} \Psi_{2} d x=1
$$

And the real parts of $\int_{-\infty}^{\infty} \Psi_{1}^{*} \Psi_{2} d x$ and $\int_{-\infty}^{\infty} \Psi_{2}^{*} \Psi_{1} d x$ are

$$
\int_{-a / 2}^{+a / 2} \cos (\pi x / a) \sin (2 \pi x / a) d x=\frac{a}{\pi} \int_{-\pi / 2}^{+\pi / 2} \cos (u) \sin (2 u) d u=0
$$

Then, in order for $\Psi$ to be normalised, it is necessary that

$$
c_{1} c_{1}^{*}+c_{2} c_{2}^{*}=1
$$

5. A particle in the infinite square well has as its initial wave function an even mixture of the first two stationary states:

$$
\Psi(x, 0)=A\left[\psi_{1}(x)+\psi_{2}(x)\right]
$$

(a) Normalise $\Psi(x, 0)$. That is, find $A$. (don't forget that $\psi(x)$ 's are orthonormal)

## solution

$$
\begin{gathered}
|\Psi|^{2}=\Psi^{*} \Psi=|A|^{2}\left(\psi_{1}^{*}+\psi_{2}^{*}\right)\left(\psi_{1}+\psi_{2}\right)=|A|^{2}\left(\psi_{1}^{*} \psi_{1}+\psi_{1}^{*} \psi_{2}+\psi_{2}^{*} \psi_{1}+\psi_{2}^{*} \psi_{2}\right) \\
1=\int|\Psi|^{2} d x=|A|^{2} \int\left(\psi_{1}^{*} \psi_{1}+\psi_{1}^{*} \psi_{2}+\psi_{2}^{*} \psi_{1}+\psi_{2}^{*} \psi_{2}\right) d x=2|A|^{2}
\end{gathered}
$$

Hence,

$$
A=\frac{1}{\sqrt{2}}
$$

(b) Find $\Psi(x, t)$ and $|\Psi(x, t)|^{2}$. Express the latter as a sinusoidal function of time using the simplification that $\omega=\pi^{2} \hbar / 2 m a^{2}$

## solution

$$
\Psi(x, t)=\frac{1}{\sqrt{2}}\left[\psi_{1} \exp \left\{-i E_{1} t / \hbar\right\}+\psi_{2} \exp \left\{-i E_{2} t / \hbar\right\}\right]
$$

Using $\psi_{n}=\sin (n \pi x / a)$ and $E_{n} / \hbar=n^{2} \omega$, this simplifies to

$$
\begin{gathered}
\frac{1}{\sqrt{a}} \exp \{-i \omega t\}\left[\sin \left(\frac{\pi}{a} x\right)+\sin \left(\frac{2 \pi}{a} x\right) \exp \{-3 i \omega t\}\right] \\
|\Psi(x, t)|^{2}=\frac{1}{a}\left[\sin ^{2}\left(\frac{\pi}{a} x\right)+\sin \left(\frac{\pi}{a} x\right) \sin \left(\frac{2 \pi}{a} x\right)(\exp \{-3 i \omega t\}+\exp \{3 i \omega t\})+\sin ^{2}\left(\frac{2 \pi}{a} x\right)\right] \\
=\frac{1}{a}\left[\sin ^{2}\left(\frac{\pi}{a} x\right)+\sin ^{2}\left(\frac{2 \pi}{a} x\right)+2 \sin \left(\frac{\pi}{a} x\right) \sin \left(\frac{2 \pi}{a} x\right) \cos (3 \omega t)\right]
\end{gathered}
$$

(c) Compute $\langle x\rangle$. Notice that it oscillates in time. What is the angular frequency of the oscillation? What is the amplitude of the oscillation?
solution

$$
\begin{aligned}
<x> & =\int x|\Psi(x, t)|^{2} d x \\
& =\frac{1}{a} \int_{0}^{a} x\left[\sin ^{2}\left(\frac{\pi}{a} x\right)+\sin ^{2}\left(\frac{2 \pi}{a} x\right)+2 \sin \left(\frac{\pi}{a} x\right) \sin \left(\frac{2 \pi}{a} x\right) \cos (3 \omega t)\right] d x
\end{aligned}
$$

First term integrate by parts

$$
\int_{0}^{a} x \sin ^{2}\left(\frac{\pi}{a} x\right) d x=\left[\frac{x^{2}}{4}-\frac{x \sin (2 \pi x / a)}{4 \pi / a}-\frac{\cos (2 \pi x / a)}{8(\pi / a)^{2}}\right]_{0}^{a}=\frac{a^{2}}{4}
$$

Second term amounts to the same value

$$
\int_{0}^{a} x \sin ^{2}\left(\frac{2 \pi}{a} x\right) d x=\frac{a^{2}}{4}
$$

Third term requires an identity to make the integral tractable

$$
\begin{aligned}
& \int_{0}^{a} \sin \left(\frac{\pi}{a} x\right) \sin \left(\frac{2 \pi}{a} x\right) d x=\frac{1}{2} \int_{0}^{a} x[\cos (\pi x / a)-\cos (3 \pi x / a)] d x \\
& =\frac{1}{2}\left[\frac{a^{2}}{\pi^{2}} \cos (\pi x / a)+\frac{a x}{\pi} \sin (\pi x / a)-\frac{a^{2}}{9 \pi^{2}} \cos (3 \pi x / a)-\frac{a x}{3 \pi} \sin (3 \pi x / a)\right]_{0}^{a} \\
& =\frac{1}{2}\left[\frac{a^{2}}{\pi^{2}}\left(\cos (\pi)-\cos (0)-\frac{a^{2}}{9 \pi^{2}}(\cos (3 \pi x / a)-\cos (0))\right]=-\frac{8 a^{2}}{9 \pi^{2}}\right.
\end{aligned}
$$

Hence the expectation value is

$$
<x>=\frac{1}{a}\left[\frac{a^{2}}{4}+\frac{a^{2}}{4}-\frac{16 a^{2}}{9 \pi^{2}} \cos (3 \omega t)\right]=\frac{a}{2}\left[1-\frac{32}{9 \pi^{2}} \cos (3 \omega t)\right]
$$

This is a function of time, with

$$
\begin{gathered}
\text { Amplitude }=\frac{32}{9 \pi^{2}} \frac{a}{2} \\
\text { angular frequency }=3 \omega=\frac{3 \pi^{2} \hbar}{2 m a^{2}}
\end{gathered}
$$

(d) If you measured the energy of this particle, what values might you get, and what is the probability of getting each of them? Find the expectation value of $H$. How does it compare with $E_{1}$ and $E_{2}$ ?
solution You could get either $E_{1}=\pi^{2} \hbar^{2} / 2 m a^{2}$ or $E_{2}=2 \pi^{2} \hbar^{2} / m a^{2}$ with equal probability $P_{1}=$ $P_{2}=1 / 2$.
So,

$$
<H>=\frac{1}{2}\left(E_{1}+E_{2}\right)=\frac{5 \pi^{2} \hbar^{2}}{4 m a^{2}}
$$

Which is the average of $E_{1}$ and $E_{2}$

